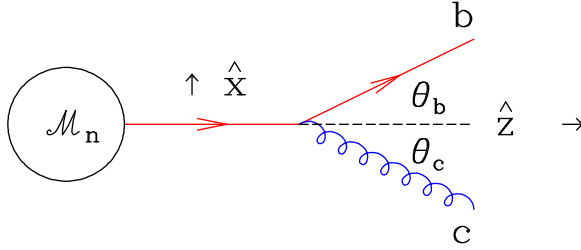


Troisième Cycle de Physique en Suisse Romande  
Exercises for Lecture 1  
Keith Ellis

1) Show that if both (massless) quarks have positive helicity the associated current for small  $\theta$  is

$$J^\mu = \bar{u}_b \gamma^\mu u_a = \sqrt{E_a E_b} (2, \theta_b, i\theta_b, 2) \quad (1)$$

with  $\theta_b$  as shown in the figure. Hence show that for small opening angle  $\theta$ , the



matrix element squared for gluon polarization in the plane is

$$|\mathcal{M}_{\text{in}}|^2 \propto \left| \frac{J_\mu \varepsilon_{\text{in}}^\mu}{t} \right|^2 = \frac{(1+z)^2}{(1-z)} \frac{1}{t} \quad (2)$$

whilst for polarization out of the plane

$$|\mathcal{M}_{\text{out}}|^2 \propto \left| \frac{J_\mu \varepsilon_{\text{out}}^\mu}{t} \right|^2 = (1-z) \frac{1}{t} \quad (3)$$

so that for

$$\frac{1}{2} [|\mathcal{M}_{\text{in}}|^2 + |\mathcal{M}_{\text{out}}|^2] \propto \frac{(1+z^2)}{(1-z)} \frac{1}{t} \quad (4)$$

2) The plus prescription is defined as

$$\int_0^1 dz f(z) g(z)_+ = \int_0^1 dz (f(z) - f(1)) g(z) \quad (5)$$

Show that

$$\int_0^1 dz z^{(n-1)} \frac{1}{(1-z)_+} = - \sum_{j=1}^{n-1} \frac{1}{j} \quad (6)$$

Using this result show that the moments of the splitting functions

$$\begin{aligned} P_{qq}(x) &= C_F \left[ \frac{1+x^2}{(1-x)_+} + \frac{3}{2} \delta(1-x) \right], \\ P_{qg}(x) &= T_R [x^2 + (1-x)^2], \quad T_R = \frac{1}{2}, \end{aligned}$$

$$\begin{aligned}
P_{gq}(x) &= C_F \left[ \frac{1 + (1-x)^2}{x} \right], \\
P_{gg}(x) &= 2C_A \left[ \frac{x}{(1-x)_+} + \frac{1-x}{x} + x(1-x) \right] \\
&\quad + \delta(1-x) \frac{(11C_A - 4n_f T_R)}{6},
\end{aligned} \tag{7}$$

are

$$\begin{aligned}
\gamma_{qq}^{(0)}(n) &= C_F \left[ -\frac{1}{2} + \frac{1}{n(n+1)} - 2 \sum_{k=2}^n \frac{1}{k} \right], \\
\gamma_{qg}^{(0)}(n) &= T_R \left[ \frac{(2+n+n^2)}{n(n+1)(n+2)} \right], \\
\gamma_{gq}^{(0)}(n) &= C_F \left[ \frac{(2+n+n^2)}{n(n^2-1)} \right], \\
\gamma_{gg}^{(0)}(n) &= 2C_A \left[ -\frac{1}{12} + \frac{1}{n(n-1)} + \frac{1}{(n+1)(n+2)} - \sum_{k=2}^n \frac{1}{k} \right] - \frac{2}{3} n_f T_R.
\end{aligned} \tag{8}$$

Hence show that for  $n = 2$  the DGLAP takes the form

$$t \frac{\partial}{\partial t} \begin{pmatrix} \Sigma(2, t) \\ g(2, t) \end{pmatrix} = \frac{\alpha_S(t)}{2\pi} \begin{pmatrix} -\frac{4}{3}C_F & \frac{1}{3}n_f \\ \frac{4}{3}C_F & -\frac{1}{3}n_f \end{pmatrix} \begin{pmatrix} \Sigma(2, t) \\ g(2, t) \end{pmatrix}. \tag{9}$$

The eigenvectors and associated eigenvalues of this system of equations are

$$\begin{aligned}
O^+(2, t) &= \Sigma(2, t) + g(2, t) \quad \text{with eigenvalue } 0, \\
O^-(2, t) &= \Sigma(2, t) - \frac{n_f}{4C_F} g(2, t) \quad \text{with eigenvalue } -\left(\frac{4}{3}C_F + \frac{n_f}{3}\right).
\end{aligned} \tag{10}$$

where

$$\Sigma(x, t) = \sum_i [q_i(x, t) + \bar{q}_i(x, t)] \tag{11}$$

and  $\Sigma(n, t)$ ,  $q(n, t)$  and  $g(n, t)$  are the corresponding moments.

$$f(n, t) = \int_0^1 dz z^{n-1} f(z) \tag{12}$$